

9/25/00

Lecture 13

Example of H bead on rotating wire

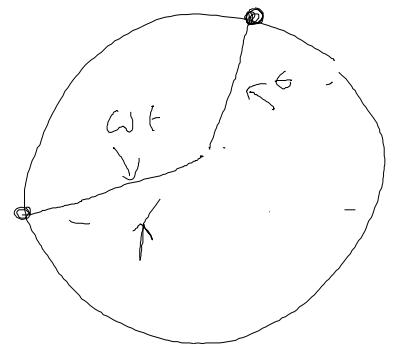
$$L = T = \frac{1}{2} m a^2 \left[\omega^2 + (\omega + \dot{\theta})^2 + 2\omega(\omega + \dot{\theta}) \cos \theta \right]$$

$$p_{\theta} = \frac{\partial L}{\partial \dot{\theta}} = m a^2 (\omega + \dot{\theta}) + m a^2 \omega \cos \theta$$

$$H = p_{\theta} \dot{\theta} - L$$

$$H = m a^2 \left[\frac{1}{2} \dot{\theta}^2 - \omega^2 (1 + \cos \theta) \right] = \text{constant}$$

$$E = T \neq \text{constant.}$$



Work problem 3.9 from problem set 3 about small osc. around stable orbit

Start reading Chapter 4 on Small Oscillations

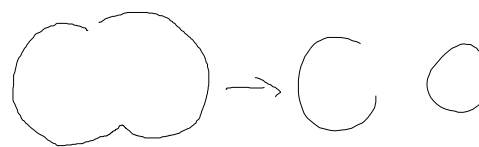
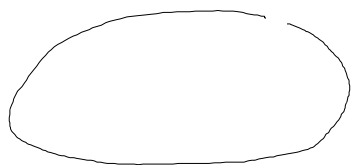
General problem to consider: stability analysis
 Do amplitude of osc. of a bridge grow with time? Is an orbit of a planet stable against small perturbations? Indeed is solar system stable? Answer is no!

We will generalize to a system of many degrees of freedom.

Also get characteristic frequencies of system. Helioseismology studies in detail spectrum of collective osc. of Sun

Examples in Nuclear physics: Is a nucleus stable against small deformations in shape?

Answer closed shell nuclei: such as ^{208}Pb
 with $N=126$ neutrons and $Z=82$ protons
 are. Many open shell nuclei: such as
 ^{235}U are unstable. Indeed true ground
 state is not spherical but deformed



Something like 30% larger on one axis.
 Note ^{235}U fissions along short axis

What are some collective osc. of nuclei
 \Rightarrow Giant dipole resonance has neutrons
 osc. against protons. Shows up
 as a large peak in photoabsorption
 cross section at $E \sim 10-20$ MeV

If you make a spherical model of ^{235}U
 discover that some mode frequencies
 are complex

with $e^{i\omega t}$ for complex ω can grow
 with time.

If you have a stable ground state
 than frequencies correspond to collective
 modes and ω^2 are all positive

Small Oscillations

Formulation (1) Look for static solutions
 of Lagrange's equations.

(2) Expand coordinates around static solution
 to lowest order in small deviations

③ Look for osc. solutions $e^{i\omega t}$
 and solve for normal mode frequencies ω .
 IF all $\omega^2 > 0$ stable, IF any $\omega^2 < 0$ than unstable.

Lagrange's equations can be written

$$\frac{d}{dt} \frac{\partial T}{\partial \dot{q}_\sigma} - \frac{\partial T}{\partial q_\sigma} = Q_\sigma = - \frac{\partial V}{\partial q_\sigma}$$

Static solution $\dot{q}_\sigma = \ddot{q}_\sigma = 0$

Generalized force vanishes $q_\sigma = q_\sigma^c$

$$Q_\sigma = - \left(\frac{\partial V}{\partial q_\sigma} \right)_{q_\sigma^c} = 0$$

IF force did not vanish it would push particles so $\dot{q}_\sigma \neq 0$

