

11/6/00

# Lecture 31 Strings and Sturm-Liouville Theory

Start reading chapter 7

Use simple physics of strings to illustrate some general math methods

Lagrangian for a string

$$L = \int_0^l \mathcal{L} dx$$

$$\mathcal{L} = \frac{1}{2} \sigma(x) \left( \frac{\partial u}{\partial t} \right)^2 - \frac{1}{2} \tau(x) \left( \frac{\partial u}{\partial x} \right)^2$$

$\sigma(x)$  = mass density

$\tau(x)$  = tension

Euler-Lagrange's eq.

$$\frac{\partial}{\partial t} \frac{\partial \mathcal{L}}{\partial \dot{u}} + \frac{\partial}{\partial x} \frac{\partial \mathcal{L}}{\partial u_x} - \frac{\partial \mathcal{L}}{\partial u} = 0$$

$$\sigma(x) \frac{\partial^2 u}{\partial t^2} = \frac{\partial}{\partial x} \left[ \tau(x) \frac{\partial u}{\partial x} \right]$$

If  $\sigma(x) = \sigma$   $\tau(x) = \tau$  a constant.

$$\frac{1}{c^2} \frac{\partial^2 u}{\partial t^2} = \frac{\partial^2 u}{\partial x^2} \quad c = \sqrt{\tau/\sigma}$$

Seek normal mode solutions

$$u(x,t) = C f(x) \cos(\omega t + \phi)$$

$$\frac{d^2 f}{dx^2} + k^2 f(x) = 0, \quad k \equiv \omega/c$$

B.C.  $p(0) = p(l) = 0$

$$p_n(x) = \left(\frac{2}{l\sigma}\right)^{1/2} \sin \frac{n\pi x}{l}$$

$$k_n = \frac{n\pi}{l} \quad n = 1, 2, 3, \dots \infty$$

$$\int_0^l dx p_m(x) \sigma p_n(x) = \delta_{nm}$$

General solution

$$u(x,t) = \sum_{n=1}^{\infty} p_n(x) \underbrace{C_n \cos(\omega_n t + \phi_n)}$$

or

$$\left[ a_n \cos\left(\frac{n\pi c t}{l}\right) + b_n \sin\left(\frac{n\pi c t}{l}\right) \right]$$

Initial conditions

$$u(x,0) = f(x)$$

$$\dot{u}(x,0) = g(x)$$

$$a_n = \left(\frac{2}{l\sigma}\right)^{1/2} \int_0^l \sin\left(\frac{n\pi x}{l}\right) f(x) \sigma dx$$

F.c.m  $\int_0^l p_n(x) \sigma u(x,0) dx$

using  $\int_0^l p_n \sigma p_m dx = \delta_{nm}$

$$\left(\frac{n\pi c}{l}\right) b_n = \left(\frac{2}{l\sigma}\right)^{1/2} \int_0^l \sin\left(\frac{n\pi x}{l}\right) g(x) \sigma dx$$

Note normal coordinates

$$y_n = c_n \cos(\omega_n t + \phi_n)$$

$$a_n = c_n \cos \phi_n \quad b_n = -c_n \sin \phi_n$$

$$L = \frac{1}{2} \sum_n (\dot{y}_n^2 - \omega_n^2 y_n^2)$$

$$H = \frac{1}{2} \sum_n (\dot{y}_n^2 + \omega_n^2 y_n^2)$$

$$H = T + V = \frac{1}{2} \sigma \int_0^l dx \left( \frac{\partial U}{\partial t} \right)^2 + \frac{1}{2} \tau \int_0^l dx \left( \frac{\partial U}{\partial x} \right)^2$$

gives same answer

$$H = E = \frac{1}{2} \sum_n \omega_n^2 c_n^2$$

### General String(like) Problem

$$\mathcal{L} = \mathcal{K} - \mathcal{V}$$

$$\mathcal{K} = \frac{1}{2} \sigma(x) \left( \frac{\partial U}{\partial t} \right)^2$$

$$\mathcal{V} = \frac{1}{2} \tau(x) \left( \frac{\partial U}{\partial x} \right)^2 + \frac{1}{2} V(x) U^2$$

Example of extra  $V(x)$  would be massive string under gravity

$$\sigma(x) \frac{\partial^2 U}{\partial t^2} = \frac{\partial}{\partial x} \left[ \tau(x) \frac{\partial U}{\partial x} \right] - V(x) U(x)$$

Seek normal modes  $U(x,t) = \rho(x) \cos(\omega t + \phi)$

Problem Set 9

Due 11/13/00

① F+W 7.3

② F+W 7.8

③ F+W 7.9

$$\boxed{-\frac{d}{dx} \left[ \tau(x) \frac{d\rho}{dx} \right] + v(x) \rho = \omega^2 \sigma(x) \rho}$$

known as Sturm - Liouville equation

$\tau(x)$ ,  $v(x)$ ,  $\sigma(x)$  are real in  $a \leq x \leq b$

$\begin{matrix} \tau(x) > 0 \\ \sigma(x) > 0 \end{matrix} \}$  in open interval  $a < x < b$

$\tau$  and  $\sigma$  could vanish at end points

B.C.

① Fixed ends  $\rho = 0$

② Natural boundary condition

③ General homogeneous b.c.  $\tau \frac{d\rho}{dx} = 0$

$$\alpha \frac{d\rho}{dx} - \beta \rho = 0$$

④ Periodic b.c.

$$\rho(a) = \rho(b)$$

$$\rho'(a) = \rho'(b)$$

Variational principle

$$\omega^2[\rho] = \frac{\frac{1}{2} \int_a^b \left[ \tau(x) \left( \frac{d\rho}{dx} \right)^2 + v \rho^2 \right] dx}{\frac{1}{2} \int_a^b \left[ \sigma(x) \rho^2 \right] dx} = \frac{I_1}{I_2} \quad 4$$

Sturm - Liouville eq. is Euler Lagrange eq. for above variational functional

$$\delta \omega^2 = \frac{1}{I_2} (\delta I_1 - \omega^2 \delta I_2) = 0$$

$$\delta \omega^2 = \frac{\delta I_1}{I_2} - \frac{I_1 \delta I_2}{I_2^2}$$

$$\delta I_1 - \omega^2 \delta I_2 = \int_a^b \left[ \tau(x) \frac{dp}{dx} \frac{d}{dx} \delta p + V p \delta p - \omega^2 \delta p \delta p \right] dx = 0$$

Integrate by parts

$$= \delta p \tau \frac{dp}{dx} \Big|_a^b + \int_a^b \left[ V p - \omega^2 \delta p - \frac{d}{dx} \left[ \tau \frac{dp}{dx} \right] \right] \delta p dx$$

⇒

$$\boxed{- \frac{d}{dx} \left[ \tau \frac{dp}{dx} \right] + V(x) p(x) = \omega^2 \delta(x) p(x)}$$

Sturm - Liouville eq.

$$IF \quad \left[ \delta p \tau p' \right] \Big|_a^b = 0$$

True for (1) fixed ends  $\delta p = 0$  (2) Natural b.c.

$\tau \rho' = 0$  with arbitrary  $Sp.$

③ Periodic b.c.

④ Need small change in functional  
for mixed b.c. see text.