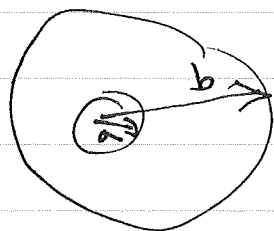


9/28/18

# lec. 9 Boundary Value Problems with Green Functions

Last time  $G$  for spherical coord.  
with  $G=0$  for  $x, x' = a$  or  $b$



$$G(\vec{x}, \vec{x}') = 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{Y_{lm}^*(\theta, \phi') Y_{lm}(\theta, \phi)}{(2l+1) \left[1 - \left(\frac{a}{b}\right)^{2l+1}\right]} \left( \frac{r^l}{r'^{l+1}} - \frac{a^{2l+1}}{r'^{2l+1}} \right) \left( \frac{1}{r'^{l+1}} - \frac{r^l}{b^{2l+1}} \right)$$

$$\Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int_V \rho(\vec{x}') G(\vec{x}, \vec{x}') d^3x' - \frac{1}{4\pi} \int_S \Phi(\vec{x}') \frac{\partial G}{\partial n'} da'$$

Consider potential inside sphere of radius  $b$ , (let  $a \rightarrow 0$ )

$$\left. \frac{\partial G}{\partial n'} \right|_{r'=b} = -\frac{4\pi}{b^2} \sum_{lm} Y_{lm}^*(\theta, \phi') Y_{lm}(\theta, \phi) \left( \frac{r'}{b} \right)^l$$

Now find  $\Phi$  inside sphere given  $\Phi(\vec{x}) = V(\theta, \phi)$  on sphere.

$$\bar{\Phi}(r, \theta, \phi) = \sum_{lm} \int_S \frac{b^2}{b^2} V(\theta', \phi') Y_{lm}^*(\theta', \phi') d\Omega' \left(\frac{r}{b}\right)^l Y_{lm}(\theta, \phi)$$

$$da' = b^2 d\Omega'$$

$$= \sum_{lm} \left[ \int_0^\pi \sin\theta' d\theta' \int_0^{2\pi} d\phi' V(\theta', \phi') Y_{lm}^*(\theta', \phi') \right] \left(\frac{r}{b}\right)^l Y_{lm}(\theta, \phi)$$

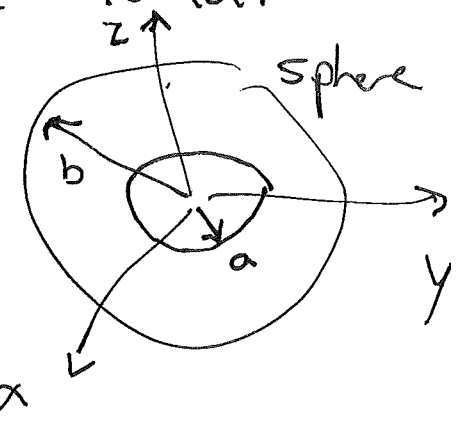
Now consider a nonzero charge density inside volume V. Enough to consider  $\bar{\Phi} \equiv 0$  on surface since we can just add above solution to satisfy  $\bar{\Phi} = V(\theta', \phi')$

Consider hollow grounded sphere of radius b with a concentric ring of radius a and total charge Q

$$\rho(\vec{x}') = \frac{Q}{2\pi a^2} \delta(r'-a) \delta(\cos\theta')$$

indep. of  $\phi'$

ring in x y plane



$$\bar{\Phi}(\vec{x}) = \frac{1}{4\pi\epsilon_0} \left(\frac{Q}{2\pi a^2}\right) \int_0^\infty r'^2 dr' \int d\cos\theta' \int d\phi' \delta(r'-a) 4\pi \sum_{lm} \frac{Y_{lm}^*(\theta', \phi') Y_{lm}(\theta, \phi)}{2l+1} r^l \left(\frac{1}{r^l} - \frac{r^l}{b^{2l+1}}\right)$$

let  $r_{<} = \text{Min}(a, r)$   
 $r_{>} = \text{Max}(a, r)$

Sdφ'  $Y_{lm}(\theta'=0, \phi) = 2\pi \delta_{m0} Y_{l0}(\theta'=0, \phi)$

$Y_{l0} = \sqrt{\frac{2l+1}{4\pi}} P_l(0)$

$P_{anti}(0) = 0$ ,  $P_{an}(0) = \frac{(-1)^n (2n-1)!!}{2^n n!}$

Note  $P_l(1) = 1$

$\Phi(\vec{x}) = \frac{Q}{8\pi^2 \epsilon_0} \sum_{n=0}^{\infty} \frac{4\pi}{(2n+1)} \sqrt{\frac{2n+1}{4\pi}} (-1)^n \frac{(2n-1)!!}{2^n n!}$

$Y_{l0}(\theta, \phi) r_{<}^{an} \left( \frac{1}{r_{>}^{2n+1}} - \frac{r_{>}^{2n}}{b^{2n+1}} \right)$

$Y_{an0} = \sqrt{\frac{2n+1}{4\pi}} P_{an}(\cos\theta)$

$\Phi(\vec{x}) = \frac{Q}{4\pi\epsilon_0} \sum_{n=0}^{\infty} (-1)^n \frac{(2n-1)!!}{2^n n!} r_{<}^{2n} \left( \frac{1}{r_{>}^{2n+1}} - \frac{r_{>}^{2n}}{b^{2n+1}} \right) P_{an}(\cos\theta)$

Now consider a line charge of total charge  $Q$  along  $z$  axis

$$\rho(\vec{x}') = \frac{Q}{2b} \frac{1}{2\pi r'^2} [\delta(\cos\theta - 1) + \delta(\cos\theta + 1)]$$

$$\Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \left(\frac{Q}{2b}\right) \frac{1}{2\pi} \int_0^b \frac{r'^2 dr'}{r'^2} \int_0^{2\pi} d\phi' \int_{-1}^1 d\cos\theta'$$

$$\sum_{lm} \frac{4\pi}{2l+1} [\delta(\cos\theta - 1) + \delta(\cos\theta + 1)] Y_{lm}^*(\theta', \phi') Y_{lm}(\theta, \phi)$$

$$r < r' \left( \frac{1}{r^{2l+1}} - \frac{r^l}{b^{2l+1}} \right)$$

$$Y_{lm}^*(\theta=0, \phi') = \delta_{m0} \sqrt{\frac{2l+1}{4\pi}} P_l(\cos\theta)$$

$$\Phi(\vec{x}) = \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{2b}\right) \int_0^b dr' \sum_l [P_l(1) + P_l(-1)] P_l(\cos\theta)$$

$$r > r' \left( \frac{1}{r^{2l+1}} - \frac{r^l}{b^{2l+1}} \right)$$

$$\int_0^b = \left( \frac{1}{r^{2l+1}} - \frac{r^l}{b^{2l+1}} \right) \int_0^r dr' r'^{-l} + r \int_r^b dr' \left( \frac{1}{r^{2l+1}} - \frac{r^l}{b^{2l+1}} \right)$$

$$= \frac{2l+1}{l(l+1)} \left[ 1 - \left(\frac{r}{b}\right)^{2l} \right] = \ln \frac{b}{r} \quad \text{for } l=0$$

$$P_\ell(-1) = (-1)^\ell$$

$$\Phi(x) = \left( \frac{Q}{4\pi\epsilon_0 b} \right) \left\{ \ln \frac{b}{r} + \sum_{j=1}^{\infty} \frac{4j+1}{(a_j)(a_j+1)} \left[ 1 - \left( \frac{r}{b} \right)^{2j} \right] P_{2j}(\cos\theta) \right\}$$

Note pot. diverges along z axis

Surface charge density

$$\sigma(\theta) = \epsilon_0 \left. \frac{\partial \Phi}{\partial r} \right|_{r=b} = -\frac{Q}{4\pi b^2} \left[ 1 + \sum_{j=1}^{\infty} \frac{4j+1}{a_j+1} P_{2j}(\cos\theta) \right]$$

$$\text{Integrate } 2\pi \int_{-1}^1 \int_{r=b}^{\infty} \sigma(\theta) r^2 dr d\cos\theta = -\frac{Q}{4\pi} 4\pi = -Q$$

Expansion of  $G$  in cylindrical coord.

$$\nabla_x^2 G(x, x') = -\frac{4\pi}{\rho} \delta(\rho - \rho') \delta(\phi - \phi') \delta(z - z')$$

$$\nabla_x^2 = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2}{\partial \phi^2} + \frac{\partial^2}{\partial z^2}$$

Write Delta Functions

$$\delta(z-z') = \frac{1}{2\pi} \int_{-\infty}^{\infty} dk e^{ik(z-z')} = \frac{1}{\pi} \int_0^{\infty} dk \cos[k(z-z')]$$

$$\delta(\phi-\phi') = \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} e^{im(\phi-\phi')}$$

Expand G

$$G(x, x') = \frac{1}{2\pi^2} \sum_{m=-\infty}^{\infty} \int_0^{\infty} dk e^{im(\phi-\phi')} \cos[k(z-z')] g_m(k, \rho, \rho')$$

$$\frac{1}{\rho} \frac{d}{d\rho} \left( \rho \frac{dg_m}{d\rho} \right) - \left( k^2 + \frac{m^2}{\rho^2} \right) g_m = -\frac{4\pi}{\rho} \delta(\rho-\rho')$$

For  $\rho \neq \rho'$  have

$$\frac{1}{\rho} \frac{d}{d\rho} \left( \rho \frac{dg_m}{d\rho} \right) - \left( k^2 + \frac{m^2}{\rho^2} \right) g_m = 0$$

let  $x' = k\rho$

$$\left[ \frac{1}{x'} \frac{d}{dx'} \left( x' \frac{dg}{dx'} \right) - \left( 1 + \frac{m^2}{x'^2} \right) \right] g = 0$$

Bessel's eq.

$$\frac{d^2 g}{dx^2} + \frac{1}{x} \frac{dg}{dx} + \left( 1 - \frac{m^2}{x^2} \right) g = 0$$

let  $x' = ix$

$$\left\{ -\frac{\partial^2}{\partial x^2} - \frac{1}{x} \frac{\partial}{\partial x} - \left(1 - \frac{m^2}{x^2}\right) \right\} g = 0$$

or

$$\frac{\partial^2 g}{\partial x^2} + \frac{1}{x} \frac{\partial g}{\partial x} + \left(1 - \frac{m^2}{x^2}\right) g = 0$$

So solutions are Bessel functions of imaginary arguments

$$I_\nu(x) \equiv i^{-\nu} J_\nu(ix)$$

$$K_\nu(x) \equiv \frac{\pi^{1/2}}{2} i^{\nu+1/2} H_\nu(ix)$$

$I_\nu$  and  $K_\nu$  are real for real  $x$  and are called ~~imaginary~~ modified Bessel functions.

If  $J_\nu$  and  $N_\nu$  are osc. then  $I_\nu$  and  $K_\nu$  are exp.

Remember for Laplace eq. can only have two osc. solutions in at most two directions  $\Rightarrow$  and exp. in third direction.

$$I_\nu \rightarrow \frac{1}{\Gamma(\nu+1)} \left(\frac{x}{a}\right)^\nu \quad x \ll 1$$

$$I_\nu \rightarrow \frac{1}{\sqrt{a\pi x}} e^x \left[1 + O\left(\frac{1}{x}\right)\right] \quad x \gg 1$$

Compare  $K_\nu \rightarrow \frac{\Gamma(\frac{\nu}{2})}{2} \left(\frac{x}{a}\right)^{\frac{\nu}{2}} \quad \nu \neq 0 \quad x \ll 1$  with  $\sqrt{\frac{\pi}{2x}} e^{-x} \quad \nu \neq 0 \quad x \gg 1$

$$J_\nu(x) \rightarrow \begin{cases} \frac{1}{\Gamma(\nu+1)} \left(\frac{x}{a}\right)^\nu & x \ll 1 \\ \sqrt{\frac{2}{\pi x}} \cos\left[x - \frac{\nu\pi}{2} - \frac{\pi}{4}\right] & x \gg 1 \end{cases}$$

Follows from  $I_\nu(x) = i^{-\nu} J_\nu(ix)$  end 9/28/10

So general solution for  $g_m(k, p, p')$   $p \neq p'$

$$g_m = \begin{cases} A(p) I_\nu(kp) + B(p') K_\nu(kp) \\ A'(p') I_\nu(kp) + B'(p') K_\nu(kp) \end{cases} \begin{matrix} p > p' \\ p < p' \end{matrix}$$

but must be symmetric in  $p, p'$   
write

$$g_m = \psi_1(kp_<) \psi_2(kp_>)$$

Normalization of product  $\psi_1, \psi_2$  determined