

## Lecture #9

## Harmonic Osc. Cont.

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + V(x) \psi = E \psi$$

$$V = \frac{1}{2} k x^2$$

$$\omega = \sqrt{k/m}$$

Dimensionless variables

$$\xi = \sqrt{\frac{m\omega}{\hbar}} x$$

$$K = \frac{E}{\left(\frac{1}{2} \hbar \omega\right)}$$

$$\frac{d^2 \psi}{d\xi^2} = (\xi^2 - K) \psi(\xi)$$

Last time

(a) Pull out asymp. behavior  $e^{\pm \xi^2/2}$   
 chose  $e^{-\xi^2/2}$  to be normalizable

(b) Guess  $\psi(\xi) = h(\xi) e^{-\xi^2/2}$

$$h'' - 2\xi h' + (K-1)h = 0$$

(c) Expand  $h = \sum_i a_i \xi^i$

$$a_{i+2} = \frac{2i+1-K}{(i+1)(i+2)} a_i$$

(d) Impose b.c.  $h(\frac{x}{2})$  can't blow up  
 $\psi(\frac{x}{2})$  will not be normalizable as  $\frac{x}{2} \rightarrow \infty$  otherwise

$\Rightarrow$  Series must terminate

$$K = 2n + 1 \quad n = 0, 1, 2, \dots$$

$$E = K \frac{1}{2} \hbar \omega = (n + \frac{1}{2}) \hbar \omega$$

With this choice of  $K$   $a_{n+2}$   
 and all higher  $n \equiv 0$

Numerical solution to see how  
 b.c. determines allowed  $E$

$$\frac{d\psi}{dx} \approx \frac{1}{2h} (\psi(x+h) - \psi(x-h))$$

$$\frac{d^2\psi}{dx^2} \approx \frac{1}{h^2} [\psi(x+h) + \psi(x-h) - 2\psi(x)]$$



Calculate  $\psi$  on grid  $\psi_i \equiv \psi(ih)$

Can see formula for 2<sup>nd</sup> der. as follows

$$\frac{d^2 \psi}{dx^2} \approx \frac{1}{h} \left[ \left. \frac{d\psi}{dx} \right|_{x+\frac{1}{2}h} - \left. \frac{d\psi}{dx} \right|_{x-\frac{1}{2}h} \right]$$

$$\text{and } \left. \frac{d\psi}{dx} \right|_{x+\frac{1}{2}h} \approx \frac{1}{h} (\psi(x+h) - \psi(x))$$

$$\left. \frac{d\psi}{dx} \right|_{x-\frac{1}{2}h} \approx \frac{1}{h} [\psi(x) - \psi(x-h)]$$

$$\psi'' = (x^2 - K) \psi(x)$$

$$\frac{1}{h^2} (\psi_{i+1} + \psi_{i-1} - 2\psi_i) \approx (x_i^2 - K) \psi_i$$

Solve for  $\psi_{i+1}$   $x_i = ih$

$$\psi_{i+1} = h^2 (x_i^2 - K) \psi_i + 2\psi_i - \psi_{i-1}$$

Consider odd solution. These have

$$\psi(0) = 0 \Rightarrow \psi_0 = 0$$

guess

$$\psi_i = \epsilon$$

$\epsilon = \text{any small \#}$

go back and normalize later

$$\psi_2 = \frac{1}{h^2} (x_1^2 - K) \epsilon + 2 \epsilon - 0$$

etc. then get  $\psi_3$  from  $\psi_2$  and  $\psi_1$

- Procedure
- Guess  $K$
  - Integrate out to large  $x$
  - Repeat and adjust  $K$  so  $\psi(x) \rightarrow 0$  as  $x \rightarrow 0$  to be normalizable.

basic code HCS, has See course  
web site Folder computer programs  
in course lectures.

$$PSIP = h^2 * (x^{n^2} - K) * PSI + 2 * PSI$$

$$PSIP = \psi_{i+1}, \quad PSIM = \psi_{i-1}$$

Shooting method, adjust  $K$  till  
Wave Function behaves as  $x \rightarrow \infty$

$$E = (n + \frac{1}{2}) h \omega \quad \text{with } n \text{ odd}$$

note  $n$  even does not have  $\psi(0) = 0$

Example of power series  
 $n=2 \Rightarrow K=5$

$$a_2 = \frac{1 - 5}{1 \cdot 2} a_0$$

$$a_2 = -2 a_0$$

$$a_4 = \frac{5 - 5}{3 \cdot 4} a_2 = 0 = a_6 \dots$$

$$\psi_{2,5}(\xi) = a_0 (1 - 2\xi^2) e^{-\xi^2/2}$$

Can find  $a_0$  by normalization

Finite polynomials  $h(\xi)$  are related to Hermite polynomials  $H_n(x)$ .

$H_n(x)$  is polynomial of order  $n$  with 1st coef  $2^n$

$$H_0 = 1, \quad H_1 = 2x, \quad H_2 = 4x^2 - 2$$

$$H_3 = 8x^3 - 12x \dots$$

Note:  $H_{n+1}(x) = 2x H_n(x) - 2n H_{n-1}(x)$

$$\psi_{2,5}(\xi) = \left( \frac{a_0}{-2} \right) H_2\left(\frac{\xi}{\sqrt{2}}\right) e^{-\xi^2/2}$$

Normalization integral can be done

$$\psi_n(x) = \left( \frac{m\omega}{\pi \hbar} \right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n\left(\frac{x}{\ell}\right) e^{-x^2/2\ell^2}$$

Normalized Harmonic osc. wave function

For low  $n$  quantum osc. very different from classical. Most prob. place to find particle in ground state  $x=0$  quantum turning points classical

$$\frac{1}{2} k x_{\text{turn}}^2 = E = V(x_{\text{turn}}) + (T=0)$$

$$x_{\text{turn}} = \pm \left( \frac{2E}{k} \right)^{1/2} = \pm \sqrt{\frac{2(n+1/2)\hbar\omega}{k}}$$

Correspondance principle

In limit  $n \rightarrow \infty$  expect to recover classical behavior.

Use code Harmonic, bas to plot wave functions.