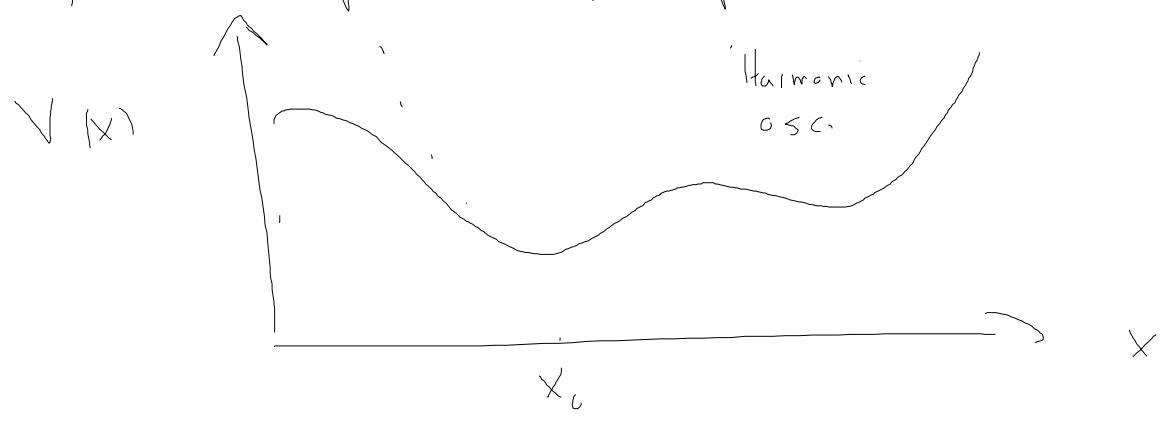


Lecture #8 Harmonic Osc.

Taylor expand any pot. about a minimum



$$V(x) \approx V(x_0) + V'(x_0)(x-x_0) + \frac{1}{2} V''(x_0)(x-x_0)^2 + \dots$$

at a minimum of energy $V(x_0)$ and $V' = 0$. Choose zero set origin so $x_0 = 0$

$$V \approx \frac{1}{2} k x^2$$

$k = \text{const.}$

Any small osc. about an equilibrium position look harmonic.

Classical freq. $\omega = \sqrt{\frac{k}{m}}$

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + \frac{1}{2} m \omega^2 x^2 \psi = E \psi$$

$$V = \frac{1}{2} k x^2 = \frac{1}{2} m \omega^2 x^2$$

Algebraic solution is "to clever"

clean up eq. a little

$\xi = \frac{m\omega}{\hbar} x$ greek xi

$$\frac{d^2\psi}{d\xi^2} = (\xi^2 - K)\psi \quad (*)$$

$$K = E / (\frac{1}{2}\hbar\omega) = \frac{2E}{\hbar\omega}$$

ξ, K are dimensionless

Solve eq. (*) by power series

(i) General equations (ugly) way to solve d.f. equations

(ii) Pull out asym. behavior first

(iii) Expand remainder in a power series

(iv) Impose b.c. to get energy eigenvalues E_n allowed values of K

(ii) For large ξ

$$\frac{d^2\psi}{d\xi^2} \approx \xi^2 \psi$$

$$\psi = A e^{-\xi^2/2} \quad B e^{+\xi^2/2}$$

bad not normalizable

(iii) Guess $\psi(\xi) = h(\xi) e^{-\xi^2/2}$

For general $h(\xi)$ this can still be exact.

$$\frac{d\psi}{d\xi} = \left[\frac{dh}{d\xi} - \xi \right] e^{-\xi^2/2}$$

$$\frac{d^2\psi}{d\xi^2} = \left[\frac{d^2h}{d\xi^2} - 2\xi \frac{dh}{d\xi} + h(\xi^2 - 1) \right] e^{-\xi^2/2}$$

So
$$\frac{d^2\psi}{d\xi^2} = (\xi^2 - K) \psi$$

$$\left[\frac{d^2h}{d\xi^2} - 2\xi \frac{dh}{d\xi} + (K - 1)h \right] = 0$$

$e^{-\xi^2/2}$ cancelled

* *

Expand

$$h = \sum_{j=0}^{\infty} a_j \xi^j$$

$$h' = \sum_{j=0}^{\infty} j a_j \xi^{j-1}$$

$$\frac{d^2h}{d\xi^2} = h'' = \sum_{j=0}^{\infty} j(j-1) a_j \xi^{j-2}$$

Change dummy variables for h''

$$h'' = \sum_{i=0}^{\infty} (i+2)(i+1) a_{i+2} \xi^i \quad i = j-2$$

$$\sum_i \left[(i+2)(i+1) a_{i+2} \xi^i - 2i a_i \xi^i + (K-1) a_i \xi^i \right]$$

If true for all ξ then each coef = 0
vanishes

$$(i+2)(i+1)a_{i+2} + \cancel{2i} [K - 2i - 1] a_i = 0$$

True for all i

$$a_{i+2} = \frac{2i+1-K}{(i+2)(i+1)} a_i$$

Recursion formula for a_{i+2} given a_i

Given a_0 get $a_2, a_4, \dots, a_{\text{even}}$

given arbitrary a_1 get $a_3, a_5, \dots, a_{\text{odd}}$

Two arbitrary constants a_0, a_1 get
 for 2nd order dif. eq. all other coef.

For large i

$$a_{i+2} \approx \frac{2}{i} a_i$$

$$a_i \approx C / \left(\frac{i}{2}\right)!$$

$$h(\xi) = C \sum_i \frac{1}{\left(\frac{i}{2}\right)!} \xi^i \approx C e^{\frac{\xi^2}{2}} = C \sum_k \frac{1}{k!} \xi^{2k}$$

If $h(\xi)$ series does not terminate will blow up as $e^{\xi^2/2}$

$$\psi(\xi) = h(\xi) e^{-\xi^2/2} \rightarrow e^{\xi^2/2}$$

gives back asympt. behavior we don't want 4

Solution to this problem is to choose some K so that there is

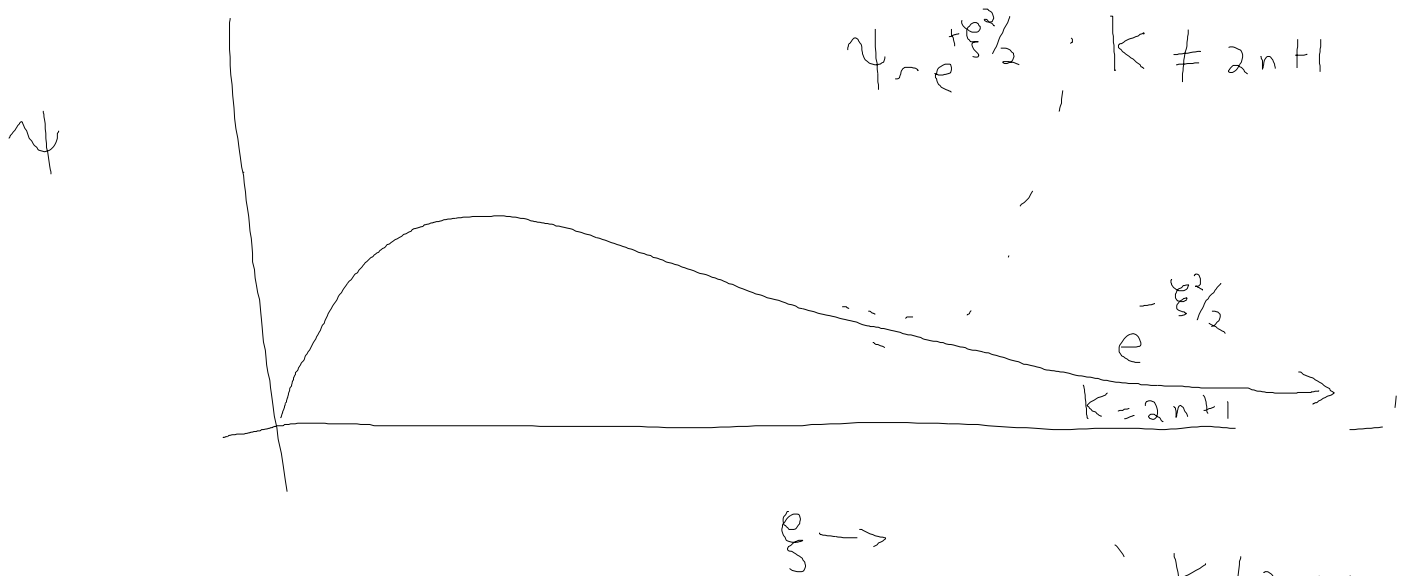
$$i = n$$

Such that

$$2n + 1 = K$$

Then $a_i \equiv 0$ for $i \geq n$

because $a_{i+2} = \frac{2i+1 - K}{(i+2)(i+1)} a_i$



Energy eigenvalues

$$K = 2n + 1$$

$$E = \frac{1}{2} \hbar \omega \quad K = (n + \frac{1}{2}) \hbar \omega$$

$n = \text{integer } 0, 1, 2, \dots$

Ground state $n=0 \Rightarrow K=1$

$$a_2 = \frac{2 \cdot 0 + 1 - 1}{(2+0)(1+0)} a_0 = 0$$

$$a_4 = \frac{3-1}{4 \cdot 3} a_2 = \frac{2}{12} a_2 = 0 = a_6 \dots$$

$$h(\xi) \equiv a_0$$

$$\psi_0 = a_0 e^{-\xi^2/2}$$

ground state

$$\xi = \left(\frac{m\omega}{\hbar}\right)^{1/2} x$$

$$\psi_0 = a_0 e^{-\left(\frac{m\omega}{\hbar}\right) \frac{x^2}{2}}$$

$$\psi_0(x) = a_0 e^{-\frac{\sqrt{m\hbar\omega}}{\hbar} \frac{x^2}{2}}$$

$$E = \frac{\hbar\omega}{2}$$

Determine a_0 by normalization

$$\int_{-\infty}^{\infty} \psi_0(x) \psi_0(x) dx = 1$$

For $n=1$ chose $a_0 \equiv c$ then $K=2n+1$

$$a_3 = \frac{[(2 \cdot 1 + 1) - (2n+1)] a_1}{3 \cdot 2} = 0 = a_5 \dots$$

$$h(\xi) = a_1 \xi$$

$$\psi_1 = a_1 \xi e^{-\xi^2/2}$$

$$E = \frac{3}{2} \hbar\omega$$

$$\psi_2 = a_0 (1 - 2\xi^2) e^{-\xi^2/2}$$

$$E = \frac{5}{2} \hbar\omega$$