

10/9/98

Lecture # 17 Matrices

Example: Pauli Spin Matrices 2×2
Hermitian matrices

$$\sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \quad \sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\sigma_y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

Properties of σ_i $i = x, y, z$

$$\sigma_i^\dagger = \sigma_i \quad \text{Hermitian}$$

Example $\sigma_y^\dagger = \sigma_y^* = \begin{bmatrix} 0 & i \\ -i & 0 \end{bmatrix}^* = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$

$$\sigma_i^2 = I \quad \sigma_y \sigma_y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

$$\left[\underline{A} \underline{B} \right]_{ik} = \sum_j A_{ij} B_{jk}$$

$$\begin{aligned} \sigma_y \sigma_y &= \begin{bmatrix} 0 \cdot 0 + (-i) \cdot i & 0 \cdot (-i) + (-i) \cdot 0 \\ i \cdot 0 + 0 \cdot i & i \cdot (-i) + 0 \cdot 0 \end{bmatrix} \\ &= \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \end{aligned}$$

Commutator

$$\begin{aligned} [\sigma_z, \sigma_x] &= \sigma_z \sigma_x - \sigma_x \sigma_z \\ \sigma_z \sigma_x &= \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} \end{aligned}$$

$$\sigma_x \sigma_z = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$$

$$\begin{aligned} [\sigma_z \sigma_x] &= \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} - \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} = +2 \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} \\ &= 2i \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} = 2i \sigma_y \end{aligned}$$

Eigenvalues and Eigenvectors

Given a transformation \hat{T} there are special vectors $|\alpha\rangle$ such that

$$\hat{T}|\alpha\rangle = \lambda|\alpha\rangle$$

Gives back a multiple λ of the original vector. λ is called the eigenvalue and $|\alpha\rangle$ an eigenvector.

Note see later Schrodinger eq. has this form

$$\hat{H}|\psi\rangle = E|\psi\rangle$$

with E the eigenvalue

Given a basis we have a Matrix eq.

$$[\underline{T}] [a] = \lambda [a]$$

or

$$\left[\underline{T} - \lambda \underline{I} \right] [a] = \underline{0}$$

Identity $n \times n$ matrix
 $n \times n$ matrix of zeros
 $\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$

If $\left[\underline{T} - \lambda \underline{I} \right]^{-1}$ exists then

$$\left[\underline{T} - \lambda \underline{I} \right]^{-1} \left[\underline{T} - \lambda \underline{I} \right] [a] = 0$$

$$\Rightarrow [a] = 0$$

So to have a nonzero eigenvector must have matrix with $\left[\underline{T} - \lambda \underline{I} \right]$ be a singular matrix with $\det = 0$.

$$\det \left[\underline{T} - \lambda \underline{I} \right] = 0$$

This is an equation for the allowed eigenvalues λ

Example: Find eigenvalues of

$$M = \begin{bmatrix} 2 & 0 & -2 \\ -2i & i & 2i \\ 1 & 0 & -1 \end{bmatrix}$$

$$|M - \lambda I| = \begin{vmatrix} 2-\lambda & 0 & -2 \\ -2i & i-\lambda & 2i \\ 1 & 0 & -1-\lambda \end{vmatrix}$$

$$= (2-\lambda)(i-\lambda)(-1-\lambda) - 0$$

$$- 2(-1)(-\lambda+i)$$

$$= -(2-\lambda)(i-\lambda)(1+\lambda) + 2(i+\lambda)$$

$$= -\lambda^3 + (1+i)\lambda^2 - i\lambda = 0$$

Det is a polynomial of order n

Roots of polynomial are 0, 1 and i

$n \times n$ matrix has at least one and at most n distinct eigenvalues.

To find eigen vector, for $\lambda = 1$
say

$$[M - I] [a] = 0$$

$$\begin{bmatrix} 1 & 0 & -2 \\ -2i & i-1 & 2i \\ 1 & 0 & -2 \end{bmatrix} \begin{bmatrix} a_1 \\ a_2 \\ a_3 \end{bmatrix} = 0$$

$$\begin{aligned} a_1 - 2a_3 &= 0 \\ -2i a_1 + (i-1)a_2 + 2i a_3 &= 0 \\ a_1 - 2a_3 &= 0 \end{aligned} \quad \text{redundant (singular matrix)}$$

let $a_1 = 1$ say $\Rightarrow a_2 = +\frac{1}{2}$

$$-2i + i(i-1)a_2 = 0$$

$$a_2 = i / (i-1) = \frac{i}{(i-1)} \frac{1+i}{1+i} = \frac{i-1}{-1-1}$$

$$a_2 = \frac{1}{2} - \frac{i}{2} = \frac{1-i}{2}$$

$$a = 1 = \begin{bmatrix} 1 \\ \frac{1}{2}(1-i) \\ \frac{1}{2} \end{bmatrix} \quad \text{for } \lambda = 1$$

Label eigen vector by its eigen value

Hermitian Transformations

(a) $T^\dagger = \hat{T}^*$

(b) $\langle T^\dagger \alpha | \beta \rangle = \langle \alpha | \hat{T} \beta \rangle$

These are equivalent

$$\langle \alpha | \beta \rangle = a^\dagger b$$

so $\langle \alpha | T \beta \rangle = a^\dagger T b = (T^\dagger a)^\dagger b$

$$\left[a^\dagger T = (T^\dagger a)^\dagger \right] \quad \longleftarrow = \langle T \alpha | \beta \rangle$$

If $T^\dagger = T$ then

① Eigenvalues of Hermitian transformation are real

calculate $T |\alpha\rangle = \lambda |\alpha\rangle$ with $|\alpha\rangle \neq 0$

$$\begin{aligned} \langle \alpha | T \alpha \rangle &= \lambda \langle \alpha | \alpha \rangle \\ &= \langle T \alpha | \alpha \rangle = \langle \lambda \alpha | \alpha \rangle \\ &= \lambda^* \langle \alpha | \alpha \rangle \end{aligned}$$

remember $\langle \alpha | \beta \rangle = \sum a_i^* b_i$ so if multiply each a_i by λ changes dot product by λ^*

Thus $\lambda \langle \alpha | \alpha \rangle = \lambda^* \langle \alpha | \alpha \rangle$ but $|\alpha\rangle \neq 0$

$$\Rightarrow \boxed{\lambda = \lambda^*}$$
 eigenvalues real.

② Eigenvalues (of Herm. trans.) are orthogonal
eigenvalues are orthogonal

$$T |\alpha\rangle = \lambda |\alpha\rangle \quad T |\beta\rangle = \mu |\beta\rangle$$

$$- \langle \alpha | T \beta \rangle = \mu \langle \alpha | \beta \rangle$$

$$- \langle T \alpha | \beta \rangle = \lambda \langle \alpha | \beta \rangle$$

Thus $\boxed{\langle \alpha | \beta \rangle = 0}$

③ Eigenvectors of a Hermitian trans. span the space

IF n distinct eigenvalues \Rightarrow n linearly indep. eigenvectors which clearly span n dim. space.

IF eigenvalue λ is m fold degenerate have m possibly non-orthogonal eigenvectors with same eigenvalue. However can orthogonalize these using Gram-Schmidt.

Note these results will be seen to imply the following features of the Harmonic osc. solutions

① Energy E is real $\hat{H}|\psi\rangle = E|\psi\rangle$

$$E = \left(n + \frac{1}{2}\right) \hbar \omega$$

(2) $\int_{-\infty}^{\infty} dx \psi_n^*(x) \psi_m(x) = 0 \quad n \neq m$

(3) Can expand any wave function

$$\bar{\Psi}(x) = \sum c_n \psi_n(x)$$

For some set ψ_n form a complete set. Thus the set.