

11/2/98

Lecture #26 H Atom

Heavy motionless proton orbited by electron of charge $-e$

$$V(r) = - \left(\frac{e^2}{4\pi\epsilon_0} \right) \frac{1}{r}$$

Full wave function

$$\Psi_{nlm}(r, \theta, \phi) = \frac{U_{nl}(r)}{r} Y_{lm}(\theta, \phi)$$

B. Conditions $U_{nl}(r \rightarrow 0) \rightarrow 0$

$U_{nl}(r \rightarrow \infty) \rightarrow 0$ so normalizable.

$$\int_0^{\infty} |U_{nl}(r)|^2 dr = 1$$

$$\int_0^{\pi} d\theta \sin\theta \int_0^{2\pi} d\phi Y_{lm}^*(\theta, \phi) Y_{lm}(\theta, \phi) = 1$$

Drop n, l label on U for now
Radial equation,

$$-\frac{\hbar^2}{2m} \frac{d^2 U}{dr^2} + \left[-\frac{e^2}{4\pi\epsilon_0} \frac{1}{r} + \frac{\hbar^2 l(l+1)}{2m r^2} \right] U = E U$$

Look for bound states $E < 0$

$$\kappa = \sqrt{\frac{-2mE}{\hbar^2}}$$

Divide by E

$$\frac{1}{\kappa^2} \frac{d^2 U}{dr^2} = \left[1 - \frac{me^2}{2\pi\epsilon_0 \hbar \kappa^2 r} + \frac{l(l+1)}{\kappa^2 r^2} \right] U$$

Define $\rho = \kappa r$, $\rho_0 = \frac{me^2}{2\pi\epsilon_0 \hbar^2 \kappa}$

$$\frac{d^2 U}{d\rho^2} = \left[1 - \frac{\rho_0}{\rho} + \frac{l(l+1)}{\rho^2} \right] U \quad (*)$$

← Coulomb pot $\sim \frac{1}{r}$
← centri. fugal term $\propto \frac{1}{r}$

- ① Power series solution both as $\rho \rightarrow \infty$ and $\rho \rightarrow 0$
- ② Pull out asympt. behavior and expand series in power series.

As $\rho \rightarrow \infty$

$$\frac{d^2 U}{d\rho^2} \approx U \Rightarrow U = A e^{-\rho} + B e^{\rho}$$

As $\rho \rightarrow 0$ centri. fugal term dominates

$$\frac{d^2 U}{d\rho^2} \approx \frac{l(l+1)}{\rho^2} U$$

$$U = C \rho^{l+1} + D \rho^{-l}$$

Guess

$$U(\rho) = \rho^{l+1} e^{-\rho} v(\rho)$$

and plug into radial eq. (*)

$$\rho \frac{d^2 v}{d\rho^2} + 2(l+1 - \rho) \frac{dv}{d\rho} + [\rho_0 - 2(l+1)] v = 0 \quad (**)$$

$$v = \sum_{j=0}^{\infty} a_j \rho^j$$

$$\frac{dv}{d\rho} = \sum_{j=0}^{\infty} (j+1) a_{j+1} \rho^j$$

$$\frac{d^2 V}{dp^2} = \sum_{j=0}^{\infty} j(j+1) a_{j+1} p^{j-1}$$

Plug into ******

$$\sum_{j=0}^{\infty} j(j+1) a_{j+1} p^j + 2(l+1) \sum_{j=0}^{\infty} (j+1) a_{j+1} p^j - 2 \sum_{j=0}^{\infty} j a_j p^j + [\rho_0 - 2(l+1)] \sum_{j=0}^{\infty} a_j p^j = 0$$

$$[j(j+1) + 2(l+1)(j+1)] a_{j+1} - 2j a_j + [\rho_0 - 2(l+1)] a_j = 0$$

$$a_{j+1} = \frac{2(j+l+1) - \rho_0}{(j+1)(j+2l+2)} a_j$$

Again argue this power series must terminate or wave function will not be normalizable.

Consider limit $j \rightarrow$ large

$$a_{j+1} \sim \frac{2}{j} a_j$$

$$\text{or } a_j \sim \frac{2^j}{j!} A$$

$$V(p) = \sum_j a_j p^j = A \sum_{j=0}^{\infty} \frac{(2p)^j}{j!} \sim e^{2p}$$

$U(\rho) = \rho^{l+1}$ $V(\rho) e^{-\rho} = \rho^{l+1} e^{-\rho}$
 and this is not normalizable
 so there must be a maximum j
 such that

$$a_{j_{\max} + 1} = 0$$

$$\text{or } 2(j_{\max} + l + 1) - \rho_0 = 0$$

Define principle quantum # n

$$n = j_{\max} + l + 1$$

$$\rho_0 = 2n = \frac{me^2}{2\pi\epsilon_0 \hbar^2} \hbar$$

$$\text{and } \hbar = \sqrt{-\frac{2mE}{\hbar^2}}$$

$$E = -\frac{\hbar^2 \hbar^2}{2m} = -\frac{me^4}{8\pi^2 \epsilon_0^2 \hbar^2 \rho_0^2}$$

$$E_n = -\frac{m}{2\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0}\right)^2 \frac{1}{n^2} = \frac{E_1}{n^2}$$

$n = 1, 2, 3, \dots$ Bohr formula

$$E_1 = -\frac{mc^2}{2} \left(\frac{e^2}{4\pi\epsilon_0 \hbar c}\right)^2$$

$$= -\frac{1}{2} mc^2 \alpha^2 = \frac{1}{2} 0.511 \text{ MeV} \left(\frac{1}{137.036}\right)^2$$

$$= -13.6 \text{ eV}$$

$$V(r) = -\frac{e^2}{4\pi\epsilon_0} \frac{1}{r} \quad \text{so } \frac{e^2}{4\pi\epsilon_0} \text{ has units energy} \cdot \text{length}$$

$$\hbar c \text{ has units energy} \cdot \text{length} = 197,33 \text{ MeV} \cdot \text{Fm}$$

$$\alpha \equiv \frac{e^2}{4\pi\epsilon_0} \frac{1}{\hbar c} \text{ is dimensionless}$$

$$\text{Fine structure constant } \alpha \approx \frac{1}{137,036}$$

Determine strength of E+M interaction

$$\lambda_C = \left(\frac{m_e c^2}{4\pi\epsilon_0 \hbar^2} \right) \frac{1}{n} = \frac{1}{a n}$$

$$a = \frac{4\pi\epsilon_0 \hbar^2}{m_e c^2} = \frac{\hbar c}{m_e c^2} \left(\frac{4\pi\epsilon_0 \hbar c}{e^2} \right) = \hbar c \left(\frac{\alpha}{m_e c^2} \right)^{-1}$$

$$= 197.33 \text{ MeV} \cdot \text{Fm} \left(\frac{1}{137.036} \right)^{-1} \left(\frac{1}{0.511 \text{ MeV}} \right)$$

$$= 5.29 \times 10^{-4} \text{ Fm}$$

$$1 \text{ Fm} = 10^{-15} \text{ m}$$

$$a = 0.529 \times 10^{-10} \text{ m}$$

Bohr radius determines size of atom

$$\rho = r / (a n)$$

$$\psi_{nlm}(r, \theta, \phi) = R_{nl}(r) Y_{lm}(\theta, \phi)$$

$$R_{nl} = \frac{1}{r} \rho^{l+1} e^{-\rho} V(\rho)$$

V is polynomial of degree

$$j_{\max} = n - l - 1 \text{ in } \rho$$

Ground state $n=1 \Rightarrow l=m=0$

$$V(\rho) = A \quad \text{const.}$$

$$\psi_{100} = A \frac{\rho^{l+1}}{r} e^{-\rho} Y_0^0(\theta, \phi)$$

$$\rho = \kappa r = r/a$$

$$\psi_{100} = \frac{A}{a} e^{-r/a} Y_0^0$$

$$Y_0^0 = \frac{1}{\sqrt{4\pi}}$$

$$\int d^3r \psi_{100}^* \psi_{100} = 1 = \int_0^\infty r^2 dr \underbrace{\int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi}_{1} \frac{(A/a)^2}{4\pi} e^{-2r/a}$$

$$= \int_0^\infty r^2 dr \frac{A^2}{a^2} e^{-2r/a} = A^2 \frac{a}{4} = 1$$

$$A = \frac{2}{\sqrt{a}}$$

$$\psi_{100} = \frac{2}{a^{3/2}} e^{-r/a} Y_0^0(\theta, \phi)$$

$$\boxed{\psi_{100}(r, \theta, \phi) = \frac{1}{\pi^{1/2} a^{3/2}} e^{-r/a}}$$