

# Lec. 25 Angular Momentum

3/19/01

## H Spectrum

$$E_\gamma = h\nu = E_i - E_f = -\frac{m_e c^2}{2} \alpha^2 \left( \frac{1}{n_i^2} - \frac{1}{n_f^2} \right)$$

$$E_i = -13.6 \text{ eV} = -\frac{m_e c^2}{2} \alpha^2$$

$$\alpha = \frac{e^2}{4\pi\epsilon_0 \hbar c}$$

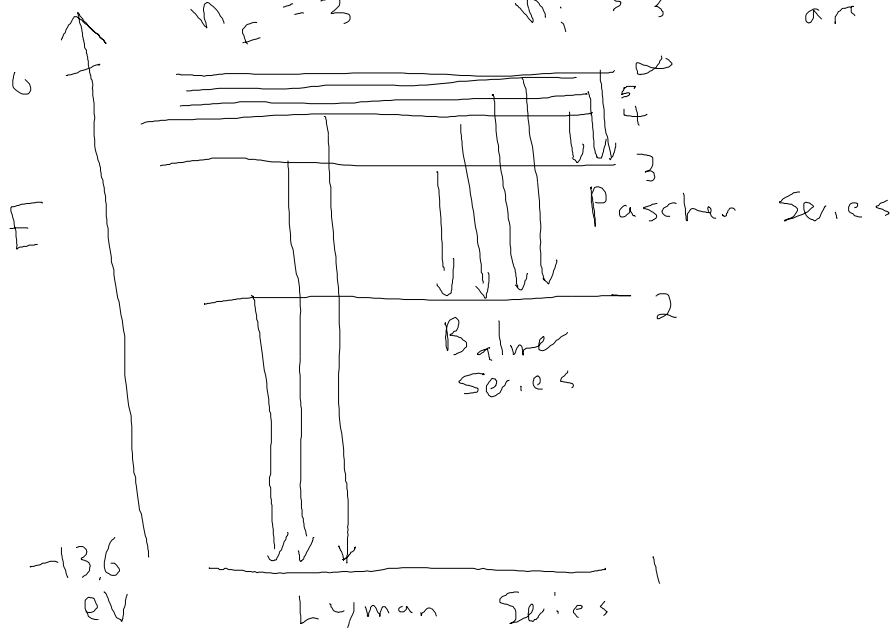
$$\nu = c/\lambda$$

$$\frac{1}{\lambda} = R \left( \frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$

$$R = \frac{m_e}{4\pi c \hbar^3} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 = 1.097 \times 10^7 \text{ m}^{-1}$$

Rydberg constant

$n_f = 1$      $n_i > 1$     are in Ultraviolet  
 (Lyman Series)  
 $n_f = 2$      $n_i > 2$     are in Visible  
 (Balmer Series)  
 $n_f = 3$      $n_i > 3$     are in Infrared  
 (Paschen Series)



## Deuterium Spectrum

Big bang produced mostly  ${}^1\text{H}$  and a small amount of  ${}^2\text{H}$  (deuterium),  ${}^3\text{He}$  and  ${}^4\text{He}$ .

Ratio of  $\text{D}/\text{H}$  from big bang provides crucial information on density of matter in Universe.

Higher  $\text{H}$  density more  $\text{D}$  gets burned into heavier  ${}^3\text{He}$  +  ${}^4\text{He}$  and lower  $\text{D}$  abundance.

How does deuterium spectrum differ from  ${}^1\text{H}$  spectrum?  $\Rightarrow$  Reduced mass effects from recoil of nucleus

Have  $\vec{r}_n = \vec{r}_{\text{nucleus}}$  and  $\vec{r}_e$  two body problem

In classical mechanics particles orbit common center of mass. Behaves like single particle with

$$\mu = \text{reduced mass} = \frac{m_e m_p}{m_e + m_p} \approx m_e$$

$$\mu_{\text{H}} \approx m_e \left(1 - \frac{m_e}{m_p}\right)$$

$$\approx m_e \left(1 - \frac{1}{2000}\right)$$

$$\mu_{\text{D}} \approx m_e \left(1 - \frac{m_e}{m_d}\right) \approx m_e \left(1 - \frac{1}{4000}\right)$$

$$\text{Since } m_d \approx m_n + m_p \approx 2m_p$$

See problem 5.1 p 178

Since  $E_i \propto \mu$

The energy of deuterium states are about one part in 4000 larger than the energy of  $^1\text{H}$  states.

$$\begin{aligned} E_i(\text{d}) &= -\frac{m_e c^2 \alpha^2}{2} \left(1 - \frac{m_e}{m_d}\right) \frac{1}{n_i^2} \\ &= -\frac{\mu c^2 \alpha^2}{2} \frac{1}{n_i^2} \end{aligned}$$

$$E_i(^1\text{H}) = -\frac{m_e c^2 \alpha^2}{2} \left(1 - \frac{m_e}{m_p}\right) \frac{1}{n_i^2}$$

Therefore deuterium spectral lines have shorter wavelengths by one part in 4000.

Example First Balmer series line has

$$\begin{aligned} \lambda^{-1} &= 1.097 \times 10^7 \text{ m}^{-1} \left(\frac{1}{2^2} - \frac{1}{3^2}\right) \\ \lambda &= 6563 \text{ \AA} \\ &\text{(in Red)} \end{aligned}$$

First Lyman Series line

$$\begin{aligned} \lambda^{-1} &= 1.097 \times 10^7 \text{ m}^{-1} \left(1 - \frac{1}{4}\right) \\ \lambda &= 1215 \text{ \AA} \quad \text{Ultraviolet} \end{aligned}$$

In a cloud at  $z = 0.701$  this light is red shifted so final wavelength is

$$\lambda = (1+z) 1215 \text{ \AA}^c$$

$$= 2066.7 \text{ \AA}$$

Correct for center of mass

$$\lambda(H) = 2066.7 / \left(1 - \frac{1}{2000}\right) = 2067.7 \text{ \AA}$$

$$\lambda(D) = 2066.7 / \left(1 - \frac{1}{4000}\right) = 2067.2 \text{ \AA}$$

$$\lambda(D) - \lambda(H) = 0.5 \text{ \AA}$$

Results

$$d / H \sim 2 \times 10^{-5}$$

Angular momentum

$$\vec{L} = \vec{r} \wedge \vec{p}$$

$$L_x = y p_z - z p_y = \frac{\hbar}{i} \left( y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right)$$

$$L_y = z p_x - x p_z = \frac{\hbar}{i} \left( z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right)$$

$$L_z = x p_y - y p_x = \frac{\hbar}{i} \left( x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$

Commutator of  $L_x$  with  $L_y$

$$[L_x, L_y] = \left(\frac{\hbar}{i}\right)^2 \left[ \left( y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \left( z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right) \right]$$

$$\begin{aligned}
& - \left( z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right) \left( y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \\
& = \left( \frac{\hbar}{i} \right)^2 \left[ y \frac{\partial}{\partial x} - x \frac{\partial}{\partial y} \right] \\
& = \frac{\hbar}{i} \left( -\frac{\hbar}{i} \right) \left[ x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right] \\
& = i\hbar L_z
\end{aligned}$$

$$\begin{aligned}
[L_x, L_y] &= i\hbar L_z \\
[L_y, L_z] &= i\hbar L_x \\
[L_z, L_x] &= i\hbar L_y
\end{aligned}$$

Can't simultaneously know  $L_x$  and  $L_y$

$$\sigma_{L_x}^2 \sigma_{L_y}^2 \geq \left( \frac{L}{2i} \langle i\hbar L_z \rangle \right)^2 = \frac{\hbar^2}{4} \langle L_z \rangle^2$$

Instead consider  $L^2$  and  $L_x$

$$[L^2, L_x] = [L_y^2, L_x] + [L_z^2, L_x]$$

$$\text{with } L^2 = L_y^2 + L_x^2 + L_z^2$$

$$\text{note } [L_x^2, L_x] = 0$$

$$\begin{aligned}
[L_y^2, L_x] &= L_y [L_y, L_x] + [L_y, L_x] L_y \\
&= -i\hbar L_y L_z - i\hbar L_z L_y
\end{aligned}$$

$$\begin{aligned}
 [L_z^2, L_x] &= L_z [L_z, L_x] + [L_z, L_x] L_z \\
 &= i\hbar L_z L_y + i\hbar L_y L_z
 \end{aligned}$$

$$\begin{aligned}
 \text{So } [L^2, L_x] &= -i\hbar L_y L_z - i\hbar L_z L_y \\
 &\quad + i\hbar L_z L_y + i\hbar L_y L_z = 0
 \end{aligned}$$

$[L^2, L_x]$	$= 0$
$[L^2, L_y]$	$= 0$
$[L^2, L_z]$	$= 0$