

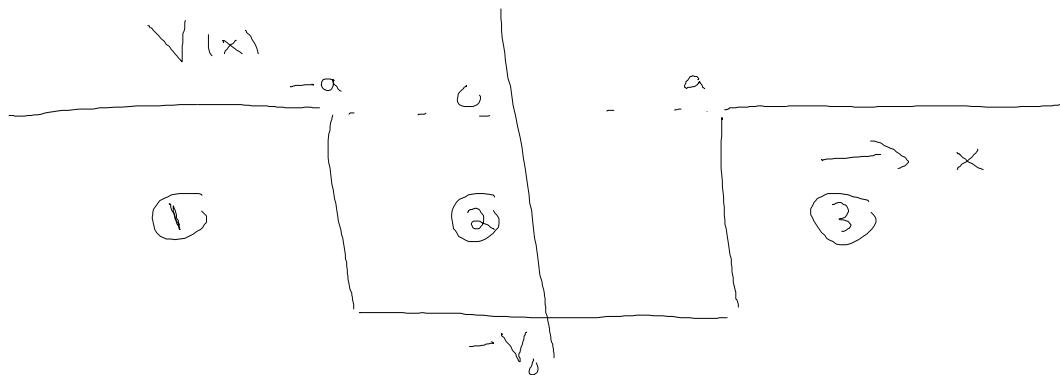
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Lecture 12 Finite Square Well

Read discussion of scattering states for a delta func. pot.

Read Sec 2.6

$$V(x) = \begin{cases} -V_0 & -a < x < a \\ 0 & \text{else} \end{cases}$$



Look for bound state $E < 0$

In region ① $V=0$ $\psi = A e^{+\kappa x}$
 so $\psi \rightarrow 0$ as $x \rightarrow -\infty$

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} = E \psi \quad \kappa = \frac{\sqrt{-2mE}}{\hbar}$$

Just like for delta function

In region ③ $V=0$ $\psi = B e^{-\kappa x}$
 $\psi \rightarrow 0$ as $x \rightarrow +\infty$

In region ② $V = -V_0$

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} - V_0 \psi = E \psi$$

$$\frac{d^2 \psi}{dx^2} = - \left[\frac{\sqrt{2m(V_0 + E)}}{\hbar} \right]^2 \psi$$

Assume $V_0 + E > 0$ even though $E < 0$

$$\text{Let } k = \frac{\sqrt{2m(E+V_0)}}{\hbar}$$

$$\psi = C \sin(kx) + D \cos kx$$

Given a solution to S. eq $\psi(x)$
Then $\psi(-x)$ is also a solution

$$\textcircled{1} \quad -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V(x) \psi(x) = E \psi(x)$$

$$\text{let } z = -x \quad dz = -dx \quad d^2z = dx^2$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(-z)}{\partial z^2} + V(-z) \psi(-z) = E \psi(-z)$$

Put is symmetric $V(-z) = V(z)$

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(-z)}{\partial z^2} + V(z) \psi(-z) = E \psi(-z)$$

$$\text{let } z \rightarrow x \quad \textcircled{2} \quad -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(-x)}{\partial x^2} + V(x) \psi(-x) = E \psi(-x)$$

So if $\psi(x)$ is a solution then $\psi(-x)$ is also a solution with same energy

Add $\textcircled{1} + \textcircled{2}$

$$\begin{aligned} -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} [\psi(x) + \psi(-x)] + V(x) [\psi(x) + \psi(-x)] \\ = E [\psi(x) + \psi(-x)] \end{aligned}$$

Now $\psi(x) + \psi(-x)$ is an even wave function.

Likewise subtract ① and ②

\Rightarrow Any solution to S. eq. for a can be taken to be even or odd. which is an even function or

Consider even solutions for square well.

$$\psi = \begin{cases} A e^{-\kappa |x|} & |x| > a \\ D \cos lx & |x| < a \end{cases}$$

B.C. (1) ~~derivative~~ wave function is cont.

(2) $\partial\psi/\partial x$ is cont. at $x=a$

$$\textcircled{1} \quad A e^{-\kappa a} = D \cos la$$

$$\textcircled{2} \quad -\kappa A e^{-\kappa a} = -l D \sin la$$

Divide ② by ①

$$\boxed{\kappa = l \tan la}$$

$$\kappa = \frac{\sqrt{-2mE}}{\hbar} \quad l = \frac{\sqrt{2m(V_0 + E)}}{\hbar}$$

m, a, V_0 are known Trans. eq. only unknown is E . F_0, E

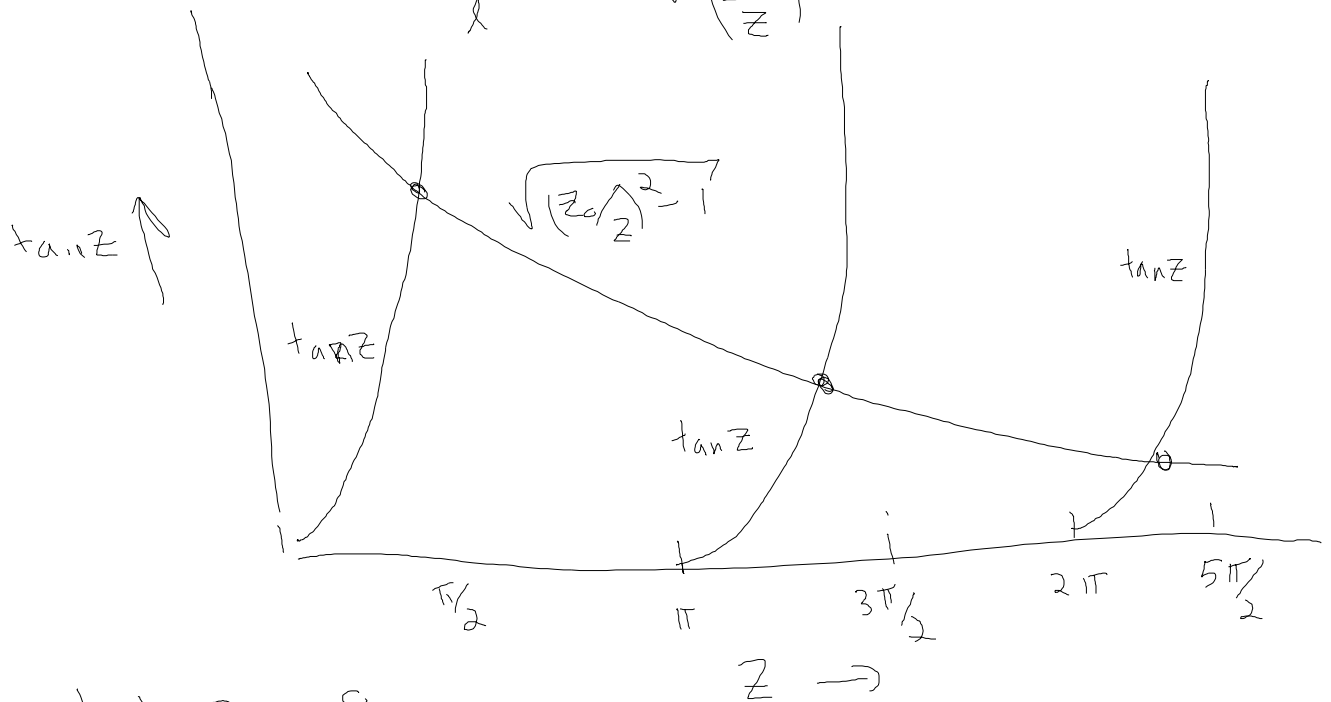
let

$$Z = ka$$

$$Z_0 = \frac{a}{\hbar} \sqrt{2mV_0}$$

$$\kappa = \frac{\sqrt{-2mE}}{\hbar}$$

$$\tan Z = \frac{\kappa}{k} = \sqrt{\left(\frac{Z_0}{Z}\right)^2 - 1}$$



let $Z_0 = 0$

Note

$$k = Z/a$$

$$\kappa^2 = \frac{-2mE}{\hbar^2}$$

$$Z^2/a^2 = k^2 = \frac{2m(V_0 + E)}{\hbar^2}$$

$$Z_0^2 = \frac{a^2}{\hbar^2} 2mV_0$$

$$Z_0^2 - Z^2 = a^2 \left[\frac{2mV_0}{\hbar^2} - \frac{2m(V_0 + E)}{\hbar^2} \right] = \frac{a^2(-2mE)}{\hbar^2}$$

so

$$\left(\frac{Z_0^2}{Z^2} - 1\right)^{1/2} = \frac{(Z_0^2 - Z^2)^{1/2}}{Z} = \frac{a}{Z} \frac{(-2mE)^{1/2}}{\hbar}$$

$$= \frac{\kappa}{k}$$

so $\left[\left(\frac{Z_0}{Z}\right)^2 - 1\right]^{1/2} = \tan Z$

Deep well. Consider limit $Z_0 \rightarrow \infty$

$$\tan Z \approx \infty$$

$$Z = n \frac{\pi}{2} \quad \text{with } n \text{ odd}$$

$$Z^2 = l^2 a^2 = \frac{2m(V_0 + E)}{\hbar^2} a^2$$

$$E = \frac{\hbar^2 Z^2}{2m a^2} - V_0$$

$$E_n = -V_0 + \frac{n^2 \hbar^2 \pi^2}{4 \cdot 2m a^2} \quad n \text{ odd}$$

Infinite square well result for $a_{tot} = 2a$

Shallow or narrow well

As $V_0 \rightarrow 0$ $E \rightarrow 0$ from below

$$l = \frac{\sqrt{2m(V_0 + E)}}{\hbar} \approx \frac{\sqrt{2mV_0}}{\hbar}$$

$$l a \approx \frac{\sqrt{2mV_0}}{\hbar} a$$

$$l a \approx l^2 a = \frac{2mV_0 a}{\hbar^2}$$

$$\frac{\sqrt{2mE}}{\hbar} = \frac{2mV_0 a}{\hbar^2}$$

$$-2mE = \frac{4mV_0^2 a^2}{\hbar^2}$$

$$E = - \frac{2mV_0^2 a^2}{\hbar^2} = - \frac{m \alpha^2}{2\hbar^2}$$

Note

$$\int_{-\infty}^{\infty} V(x) dx = V_0 2a = \alpha$$