

12/11/98

Lecture #40 Final Review

1) Time dependence in QM given by

$$i\hbar \frac{\partial \Psi}{\partial t} = \hat{H} \Psi$$

2) Look for stationary states

$$\Psi(x,t) = \psi(x) e^{-iEt/\hbar}$$

$$\hat{H} \psi(x) = E \psi(x)$$

3) Nontrivial time dependence only for superposition of stationary states

$$\Psi(x,t) = \sum_i c_i \psi_i(x) e^{-iE_i t/\hbar}$$

Can always expand any wave function in stationary states.

4) Solution to 1 dim $\hat{H} \psi(x) = E \psi(x)$

a) Wave function is always cont.

b) Derivative $\frac{d\psi}{dx}$ is cont. except where $V \rightarrow \infty$

c) Adjust E to satisfy boundary conditions

So that $\psi(x) \rightarrow 0$ as $x \rightarrow \pm\infty$
So that $\int dx \psi^* \psi < \infty$

5) General analytic solution of 1dim S. eq.
 Example: H. osc.

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + \frac{1}{2} m \omega^2 x^2 \psi = E \psi$$

$$\xi = \sqrt{m\omega/\hbar} x \quad K = 2E/\hbar\omega$$

$$\frac{d^2 \psi}{d\xi^2} = -(\xi^2 - K) \psi$$

a) Find large ξ solution $\frac{d^2 \psi}{d\xi^2} = \xi^2 \psi \quad \psi \propto e^{-\xi^2/2}$

$$\Rightarrow \psi = h(\xi) e^{-\xi^2/2}$$

b) Expand $h(\xi)$ in power series $h = \sum_i a_i \xi^i$

c) Adjust E so wave function is normalizable (series terminates) $E = (n + \frac{1}{2}) \hbar\omega$

Three dimensions

Separation of variables for spherical sym.
 pot.

$$V(\vec{r}) \rightarrow V(r)$$

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2}$$

$$\Psi(r, \theta, \phi) = R(r) Y_l(\theta, \phi)$$

Angular eq.

$$\sin \theta \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta} Y + \frac{\partial^2 Y}{\partial \phi^2} = -l(l+1) \sin^2 \theta Y$$

$$Y_l = Y_l^m(\theta, \phi) \quad \text{spherical Harmonic}$$

$$\int_0^\pi d\theta \sin \theta \int_0^{2\pi} d\phi Y_{l'}^{*m'} Y_l^m = \delta_{mm'} \delta_{ll'}$$

$$R(r) = U(r) / r \quad \text{reduced radial wave function}$$

Radial eq.

$$-\frac{\hbar^2}{2m} \frac{d^2}{dr^2} U_{nl} + \left[V(r) + \frac{\hbar^2}{2m} \frac{l(l+1)}{r^2} \right] U_{nl}(r) = E U_{nl}$$

$$\Psi_{nlm}(r, \theta, \phi) = \frac{U_{nl}(r)}{r} Y_l^m(\theta, \phi)$$

B.C. $U \rightarrow 0$ as $r \rightarrow 0$
 $r \rightarrow \infty$

F_0 wave function to be normalizable.

Review solution for H atom $V = \frac{-e^2}{4\pi\epsilon_0 r}$

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Angular momentum

$$L = \vec{r} \wedge \vec{p}$$

$$[L_x, L_y] = i\hbar L_z$$

but $[L_i, L^2] = 0 \quad i = x, y, z$

so can have simultaneous eigenvalues
for L^2, L_z

$$L^2 Y_l^m = \hbar^2 l(l+1) Y_l^m$$

$$L_z Y_l^m = \hbar m Y_l^m \quad m = -l, -l+1, \dots, l+1, l$$

Addition of angular momentum

$$\vec{J} = \vec{L} + \vec{S}$$

J_z has eigenvalues $m_z = m + m_s$

J^2 has eigenvalues $\hbar^2 j(j+1)$

allowed $j = |l-s|, |l-s|+1, \dots, l+s$
From triangle rule



Quantum Stat. Mechanics

a) $\frac{1}{2}$ integer spins are Fermions: Wave function must be antisymmetric under interchange of any two identical Fermions

b) Integer spins are bosons: wave function is sym.

c) Pauli exclusion principle

d) Maximize number of possible states for system in thermal equilibrium gives Fermi-Dirac dist.

$$n(E) = \left[e^{(E-\mu)/kT} + 1 \right]^{-1}$$

or Bose-Einstein dist.

$$n(E) = \left[e^{(E-\mu)/kT} - 1 \right]^{-1}$$

e) # of μ particles to get corr. l.